

physics of Fibre bundles

EM theory : a vector field on the base manifold $M \equiv \mathbb{R}^2$

$$\vec{E}(\underbrace{x, y}_P) = (xy^2, xy^3)$$

In Math : The collection of all tangent spaces of $M \equiv \mathbb{R}^2$

$$TM = \bigcup_{P \in M} T_P M$$

(total space)

$$\begin{array}{ccc} E & \xrightarrow{\pi} & M \\ = TM & \downarrow \text{projection} & \end{array}$$

A section (or a vector field) of TM is a smooth map.

$$s : M \rightarrow TM \text{ such that } \pi \circ s = \text{Id}_M \quad (s(p) = \pi^{-1}(p) = u \in TM)$$

$p \rightarrow \pi^{-1}(p)$

$$\begin{array}{ccc} M & & F \\ \downarrow & \swarrow & \\ & & \end{array}$$

For this case : $TM = \mathbb{R}^2 \times \mathbb{R}^2$

a trivial bundle is a direct product

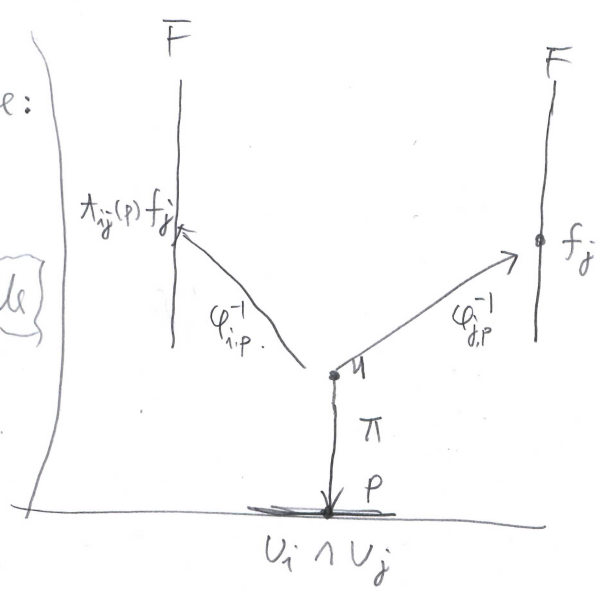
$$\begin{array}{ccc} M \times F & & \\ \uparrow & \uparrow & \\ \text{base} & & \text{fibre} \end{array}$$

A simple non-trivial fibre bundle.

$M = S^1$, $F = [-1, 1]$ = a line segment

$E \xrightarrow{\pi} S^1$ structure group = $G = Z_2 = \{e, g\}$

For the important case:
 $F \cong G$
 \Rightarrow principal fibre bundle
 \Rightarrow connections, physics.

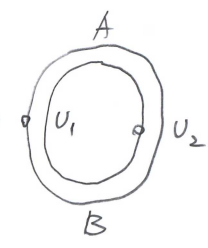


τ_{ij} : transition function
 ϕ_i : local trivialization

$\{U_i\}$: an open covering of M .

Let $U_1 = (0, 2\pi)$, $U_2 = (-\pi, \pi)$ ← a covering of S^1

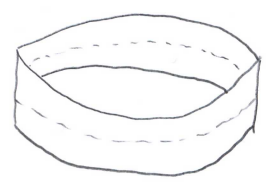
and $A = (0, \pi)$, $B = (\pi, 2\pi)$, the intersection $U_1 \cap U_2$



(I) $\phi_1^{-1}(u) = (\theta, x)$, $\theta \in A$; $\phi_2^{-1}(u) = (\theta, x)$, $\theta \in B$

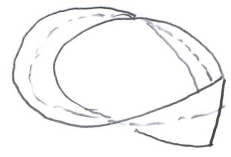
(II) $\phi_1^{-1}(u) = (\theta, x)$, $\theta \in A$; $\phi_2^{-1}(u) = (\theta, -x)$, $\theta \in B$

(I)



cylinder, trivial

(II)



Möbius strip.

Complex Manifolds

Recall: One can easily def a vector field on the plane \mathbb{R}^2 .

$\boxed{[121]}$ $\vec{E}(x, y) = (xy^2, xy^2)$ where (x, y) : coordinates of \mathbb{R}^2 .

But how can one def a vector field on S^2 ?

\otimes { A cross-section of the tangent bundle TS^2 } Martin P311

It turns out that one needs to introduce at least

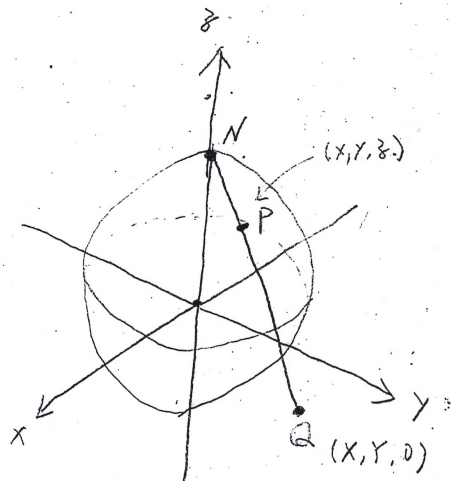
two coordinate systems, say 2 stereographic coordinates on S^2 .

• Stereographic coordinates (x, y) of S^2

$$x = \sin \theta \cos \varphi, \quad y = \sin \theta \sin \varphi, \quad z = \cos \theta$$

$$\left\{ \begin{aligned} \theta &= \tan^{-1} \left(\frac{(x^2 + y^2)^{1/2}}{z} \right), \quad \varphi = \tan^{-1} \frac{y}{x} \end{aligned} \right.$$

$$X = \frac{x}{1-z}, \quad Y = \frac{y}{1-z}$$



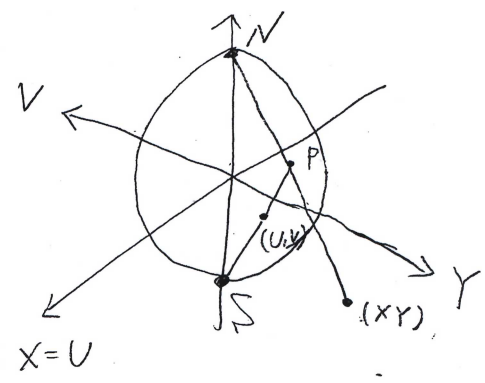
$$\Rightarrow X = \cot \frac{\theta}{2} \cos \varphi, \quad Y = \cot \frac{\theta}{2} \sin \varphi$$

Note: (θ, φ) 不好磨擦: $\textcircled{1}$ $\varphi: 0 \rightarrow 2\pi$ 不連續

(X, Y) 也不好 : N 真座標 = ??

One way out ! \exists 1 條 ≥ 2 組座標 (X, Y) & (U, V)

(X, Y) 可以覆蓋 $S^2 - \{N\}$
 (U, V) - - - $S^2 - \{S\}$



那麼 (X, Y) 跟 (U, V) 的關係是 ?

P 點可用 (X, Y) 座標表之。
or (U, V)

Ans $\begin{cases} (X, Y) = \left(\frac{x}{1-z}, \frac{y}{1-z} \right) \\ (U, V) = \left(\frac{x}{1+z}, \frac{-y}{1+z} \right) \end{cases}$: Cover S^2 with an atlas of two charts

trick : 令 $\begin{cases} z = x+iy, \bar{z} = x-iy \\ w = u+iv, \bar{w} = u-iv \end{cases}$

則 $w = \frac{x-iy}{1+z} = \frac{1-z}{1+z} (x-iy) = \frac{x-iy}{x^2+y^2} = \frac{1}{z}$

$\therefore \begin{cases} U(x, Y) = \frac{x}{x^2+y^2} \\ V(x, Y) = \frac{-y}{x^2+y^2} \end{cases}$

$U(x, Y)$ 及 $V(x, Y)$ 在赤道附近 均為 C^∞ 函數。

smoothly !

Now we can def a vector field on S^2

by using two charts (X, Y) and (U, V)

例記 $X \frac{\partial}{\partial X} + Y \frac{\partial}{\partial Y}, -U \frac{\partial}{\partial U} - V \frac{\partial}{\partial V}$

定义了 一个 vector field on S^2

proof: We need to prove that the two expressions agree on the intersection of the two coordinate systems.

$$-U \frac{\partial}{\partial U} - V \frac{\partial}{\partial V} = -\frac{X}{X^2+Y^2} \left(\frac{\partial X}{\partial U} \frac{\partial}{\partial X} + \frac{\partial Y}{\partial U} \frac{\partial}{\partial Y} \right) - \left(\frac{-Y}{X^2+Y^2} \right) \left(\frac{\partial X}{\partial V} \frac{\partial}{\partial X} + \frac{\partial Y}{\partial V} \frac{\partial}{\partial Y} \right)$$

$$\begin{aligned} \left[\frac{\partial}{\partial X} \right] &: -\frac{X}{X^2+Y^2} \frac{V^2-U^2}{[U^2+V^2]^2} + \frac{Y}{X^2+Y^2} \frac{-2UV}{[U^2+V^2]^2} \\ &= -X(X^2+Y^2)(V^2-U^2) - Y(X^2+Y^2)2UV \\ &= X^2+Y^2 \left[-X \frac{Y^2-X^2}{[X^2+Y^2]^2} + Y \frac{2XY}{[X^2+Y^2]^2} \right] \\ &= \frac{XY^2+X^3}{X^2+Y^2} = X \end{aligned}$$

$$\left[\frac{\partial}{\partial Y} \right]$$

$$\begin{aligned} & -\frac{X}{X^2+Y^2} \frac{2UV}{[U^2+V^2]^2} + \frac{Y}{X^2+Y^2} \frac{V^2-U^2}{[U^2+V^2]^2} \\ &= -X(X^2+Y^2)2UV + Y(X^2+Y^2)(V^2-U^2) \\ &= X^2+Y^2 \left[-X \frac{(-2XY)}{[X^2+Y^2]^2} + Y \frac{Y^2-X^2}{[X^2+Y^2]^2} \right] \\ &= \frac{YX^2+Y^3}{X^2+Y^2} = Y \end{aligned}$$

∴ 在赤道附近

$$-U \frac{\partial}{\partial U} - V \frac{\partial}{\partial V} = X \frac{\partial}{\partial X} + Y \frac{\partial}{\partial Y} \quad \#$$

[Note]

There are 2 zeros of the vector field.

North pole $(U, V) = (0, 0)$

South pole $(X, Y) = (0, 0)$

("hairy ball" theorem)

• A vector bundle TS^2 : the tangent bundle of S^2

An open covering of S^2 : $U_N \equiv S^2 - \{S\}$, $U_S \equiv S^2 - \{N\}$

(X, Y) be the stereographic coordinates of U_N

(U, V) - - - - - U_S

$$\Rightarrow U = \frac{X}{X^2+Y^2}, \quad V = \frac{-Y}{X^2+Y^2}$$

$$u \in TM, \quad \pi(u) = p \in U_N \cap U_S,$$

$$\varphi_N^{-1}(u) = (p, v_N^u), \quad \varphi_S^{-1}(u) = (p, v_S^u) \leftarrow \text{local trivializations}$$

the transition function is

$$T_{v_S}^{-1} = T_{v_N} = \frac{\partial(U, V)}{\partial(X, Y)} = \begin{pmatrix} \frac{\partial U}{\partial X} & \frac{\partial U}{\partial Y} \\ \frac{\partial V}{\partial X} & \frac{\partial V}{\partial Y} \end{pmatrix}$$

$$\begin{pmatrix} X = r \cos \theta \\ Y = r \sin \theta \end{pmatrix}$$

$$= r^{-2} \begin{pmatrix} -\cos \theta & -\sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix}$$

$= 4\pi$ for $\theta = 2\pi$

2\theta!!!

$$\frac{\partial U}{\partial X} = \frac{\partial r}{\partial X} \frac{\partial U}{\partial r} + \frac{\partial \theta}{\partial X} \frac{\partial U}{\partial \theta} = \cos \theta \left(-\frac{\cos \theta}{r^2}\right) + \left(-\frac{\sin \theta}{r}\right) \left(\frac{-\sin \theta}{r}\right) = \frac{-\cos 2\theta}{r^2}$$

$$\frac{\partial U}{\partial Y} = \frac{\partial r}{\partial Y} \frac{\partial U}{\partial r} + \frac{\partial \theta}{\partial Y} \frac{\partial U}{\partial \theta} = \sin \theta \left(-\frac{\cos \theta}{r^2}\right) + \left(\frac{\cos \theta}{r}\right) \left(\frac{-\sin \theta}{r}\right) = \frac{-\sin 2\theta}{r^2}$$

$$\frac{\partial V}{\partial X} = \frac{\partial r}{\partial X} \frac{\partial V}{\partial r} + \frac{\partial \theta}{\partial X} \frac{\partial V}{\partial \theta} = \cos \theta \left(\frac{\sin \theta}{r^2}\right) + \left(-\frac{\sin \theta}{r}\right) \left(\frac{-\cos \theta}{r}\right) = \frac{\sin 2\theta}{r^2}$$

$$\frac{\partial V}{\partial Y} = \frac{\partial r}{\partial Y} \frac{\partial V}{\partial r} + \frac{\partial \theta}{\partial Y} \frac{\partial V}{\partial \theta} = \sin \theta \left(\frac{\sin \theta}{r^2}\right) + \left(\frac{\cos \theta}{r}\right) \left(\frac{-\cos \theta}{r}\right) = \frac{-\cos 2\theta}{r^2}$$

• The Hopf fibration

(-) 知道 S^3

S^3 is an $U(1) \cong S^1$ bundle over S^2 . (1931)

- 先 change notation of parametrization of S^2

$$(x, y, z) \rightarrow (\xi^1, \xi^2, \xi^3), \quad (\xi^1)^2 + (\xi^2)^2 + (\xi^3)^2 = 1$$

✱

- Need parametrization of S^3 :

$$(x^1)^2 + (x^2)^2 + (x^3)^2 + (x^4)^2 = 1$$

Def: $z^0 \equiv x^1 + ix^2$, $z^1 \equiv x^3 + ix^4 \Rightarrow |z^0|^2 + |z^1|^2 = 1$

✱

- ✱ If we do the following identification

$$Z \equiv \frac{\xi^1 + i\xi^2}{1 - \xi^3} = \frac{\overbrace{x^1 + ix^2}^{S^3}}{\underbrace{x^3 + ix^4}_{S^3}} \equiv \frac{z^0}{z^1}, \quad \text{for } \xi \in U_S,$$

$$W \equiv \frac{\xi^1 - i\xi^2}{1 - \xi^3} = \frac{\overbrace{x^3 + ix^4}^{S^3}}{\underbrace{x^1 + ix^2}_{S^3}} \equiv \frac{z^1}{z^0}, \quad \text{for } \xi \in U_N,$$

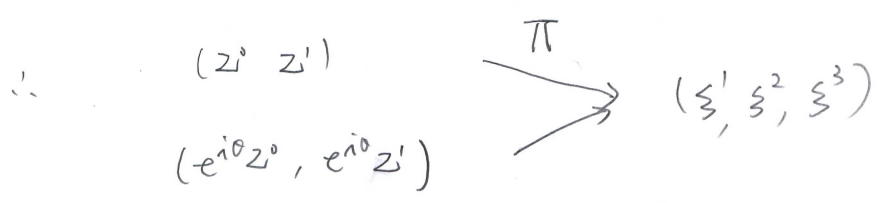
We get the Hopf map $\pi: S^3 \rightarrow S^2$

$$\begin{cases} \xi^1 = 2(x^1x^3 + x^2x^4) \\ \xi^2 = 2(x^2x^3 - x^1x^4) \\ \xi^3 = (x^1)^2 + (x^2)^2 - (x^3)^2 - (x^4)^2 \end{cases}$$

Observation

Z is invariant under $(z^0, z^1) \rightarrow (e^{i\theta} z^0, e^{i\theta} z^1)$

where $e^{i\theta} \in U(1) \Rightarrow (e^{i\theta} z^0, e^{i\theta} z^1) \in S^3$



Indeed for all (x^1, x^2, x^3, x^4) on S^3 with

$$\frac{x^1 + ix^2}{x^3 + ix^4} = \frac{x^1 + ix^2}{x^3 + ix^4} \quad (x^1, \dots, x^4 \text{ fixed})$$

$$\text{or } \begin{cases} x^3 x^1 - x^4 x^2 - x^1 x^3 + x^2 x^4 = 0 \\ x^4 x^1 + x^3 x^2 - x^2 x^3 - x^1 x^4 = 0 \end{cases} \quad (*)$$

The images under π are the same (ξ^1, ξ^2, ξ^3) .

But (*) is a great circle of S^3 .

"Fibre"
 S^1

拾遺

$$\bullet \pi : S^7 \rightarrow S^4 \quad \text{quaternions}, \quad F = SU(2) = S^3$$

$$\bullet \pi : S^{15} \rightarrow S^8 \quad \text{Octonions (Cayley numbers)}, \quad F = S^7, \text{ but NOT a group!}$$

$$\pi : S^{2n-1} \rightarrow S^n$$

$n = 2, 4, 8$

• The fibre bundle structure

(=) 不知道 S^3

local trivialisations:

$$\varphi_S^{-1} : \pi^{-1}(U_S) \rightarrow U_S \times U(1)$$

$$(z^0, z^1) \rightarrow (z^0/z^1, z^1/z^1),$$

$$\varphi_N^{-1} : \pi^{-1}(U_N) \rightarrow U_N \times U(1)$$

$$(z^0, z^1) \rightarrow (z^1/z^0, z^0/z^0).$$

On the equator $\xi^3 = 0 \Rightarrow |z^0| = |z^1| = \frac{1}{\sqrt{2}}$

∴ on the equator $\varphi_S^{-1} : (z^0, z^1) \rightarrow (z^0/z^1, \sqrt{2}z^1)$,

$$\varphi_N^{-1} : (z^0, z^1) \rightarrow (z^1/z^0, \sqrt{2}z^0).$$

The transition function on the equator is

$$A_{NS}(\xi) = \frac{\sqrt{2}z^0}{\sqrt{2}z^1} = \xi^1 + i\xi^2 \in U(1)$$

∴ Topological Charge $Q = 1$ of $\pi_1(U(1)) = \mathbb{Z}$ group.

$$S^1 \rightarrow S^1$$

- Alternative description of Hopf map (\equiv) 知道 S^3

\Rightarrow 推廣 S^1

$S^3 \leftrightarrow$ a complex 1-sphere S_c^1

$$S_c^1 = \{ (z^0, z^1) \in \mathbb{C}^2 \mid |z^0|^2 + |z^1|^2 = 1 \}$$

Def $\pi : S_c^1 \rightarrow \mathbb{C}P^1$

$$(z^0, z^1) \rightarrow \{ [z^0, z^1] \} = \{ \lambda (z^0, z^1) \mid \lambda \in \mathbb{C}, \lambda \neq 0 \}$$

points of S^3 of the form $\lambda (z^0, z^1)$, $|\lambda|=1$ are mapped to

a single point of $\mathbb{C}P^1 \cong S^2$.

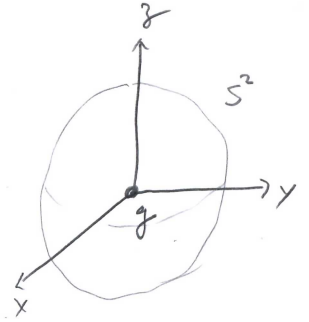
- Dirac monopole (1931 !! same year with Hopf)

(10) 解物理 \vec{A}

physical realization of Hopf!

- Maxwell theory is a U(1) Gauge Theory.

- To avoid Dirac string, one needs to use 2-patches to describe monopole fields:



g: magnetic monopole.

$$1. \quad A_x^N = \frac{-g y}{r(r+z)}, \quad A_y^N = \frac{g x}{r(r+z)}, \quad A_z^N = 0.$$

\Rightarrow one has $\nabla \times \vec{A} = \vec{B}$ except along the negative-z axis.

$$2. \quad A_x^S = \frac{g y}{r(r-z)}, \quad A_y^S = \frac{-g x}{r(r-z)}, \quad A_z^S = 0$$

\Rightarrow one has $\nabla \times \vec{A} = \vec{B}$ except along the positive-z axis.

(HW) In polar coordinates

$$\left\{ \begin{array}{l} \vec{A}^N(\vec{r}) = \frac{g(1-\cos\theta)}{r\sin\theta} \hat{e}_\varphi \\ \vec{A}^S(\vec{r}) = -\frac{g(1+\cos\theta)}{r\sin\theta} \hat{e}_\varphi \end{array} \right., \quad \hat{e}_\varphi = -\sin\theta \hat{e}_x + \cos\theta \hat{e}_y.$$

$$\vec{A}^N - \vec{A}^S = \nabla \overbrace{2\theta\phi}^{\equiv \Lambda} \quad (\text{on equator})$$

G.T. $\left\{ \begin{array}{l} \psi^S(\vec{r}) = e^{\frac{-ie\Lambda}{\hbar c}} \psi^N(\vec{r}) \end{array} \right. \leftarrow \text{electron wave function} \\ \text{(a section)}$

wave function single-valued

$$\Rightarrow \frac{e}{\hbar c} (2g \cdot \underbrace{2\pi}_{\Delta\phi}) = 2\pi n$$

$$\Rightarrow \boxed{\frac{2eg}{\hbar c} = n} \leftarrow \text{整数}$$

Dirac quantization condition $\Rightarrow \otimes e$ quantized !!!

$n \equiv 1$ for magnetic monopole.

Monopole bundle \Leftrightarrow $U(1)$ principal bundle over S^2
 $\equiv S^1$

Dirac (1931)

Hopf (1931)

Wu-Yang (1968)

Hopf map $S^7 \rightarrow S^4$

(\equiv) $\mathbb{R}P^1 \times S^7$

quaternion

$$i^2 = j^2 = k^2 = -1$$

$$ij = -ji = k$$

$$jk = -kj = i$$

$$ki = -ik = j$$

$$\left\{ \begin{array}{l} q = x + xi + yj + zk \\ |q| = (x^2 + x^2 + y^2 + z^2) \end{array} \right.$$

$$|q| = 1 \quad \text{表} \quad SU(2) = S^3$$

S^7 is a quaternion one-sphere S_H^1

$$S_H^1 = \{ (q^0, q^1) \in H^2 \mid |q^0|^2 + |q^1|^2 = 1 \}$$

$$\text{Def } \pi : S_H^1 \rightarrow HP^1$$

$$(q^0, q^1) \rightarrow \{ [q^0, q^1] \} = \{ \eta (q^0, q^1) \mid \eta \in H, \eta \neq 0 \}$$

The transition function defined by this Hopf map

belongs to the class 1 of $\pi_3(SU(2)) = \mathbb{Z}$ \leftarrow 1-instanton

points of S^7 of the form $\eta (q^0, q^1)$, $|\eta|=1$ are

mapped to a single point of $HP^1 \equiv S^4$.

• The fibre bundle structure of instanton $(=) \approx \mathbb{R}^4 \times S^3$

$$\left\{ \begin{aligned} U_N &= \{ (x, y, z, t) \mid x^2 + y^2 + z^2 + t^2 \leq R^2 + \epsilon \} \\ U_S &= \{ (x, y, z, t) \mid x^2 + y^2 + z^2 + t^2 \geq R^2 - \epsilon \} \end{aligned} \right.$$

at a point $P = (x, y, z, t) \in U_N \cap U_S$, $\pi^{-1}(P) = \mathcal{U}$.

The local trivializations.

$$\varphi_N^{-1}(\mathcal{U}) = (P, g_N), \quad \varphi_S^{-1}(\mathcal{U}) = (P, g_S)$$

where $g_N, g_S \in SU(2)$

on $U_N \cap U_S$, we have $g_N = \frac{1}{R} (tI + i x^i \sigma_i) g_S$

the large gauge transformation.

where $\frac{tI + i x^i \sigma_i}{\sqrt{x^2 + y^2 + z^2 + t^2}} (\in \mathbb{R})$ is the transition function.
(on large gauge transf)

$$Q = 1 \quad \text{of} \quad \pi_3(SU(2)) = \mathbb{Z} \text{ group.}$$

1-Instanton bundle(10) 解 物理 SDYM

$$A'_m(x) = \eta_{mL} \frac{(x-a)_L}{|x-a|^2 + \lambda^2}$$

↙ attach symbol

$$F'_{mL}(x) = \eta_{mL} \frac{\lambda^2}{[|x-a|^2 + \lambda^2]^2}$$

$$Q = 1$$